continuously. Moreover, from the MZ calculation, $T_K$ would then have to attain a value $\sim 10^6$ °K to account for the observed reduction of $\Delta T_c$ at pressures $\gtrsim 100$ kbar, which seems quite unphysical. So far in our discussion we have neglected the possibility of anomalous behavior of $N(0)$ under pressure, which cannot be completely ruled out.

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[16] R. M. More [Ph.D. thesis, University of California, San Diego, 1968 (unpublished)] has simplified the result of Suhl and Wong (Ref. 8) by replacing the finite range interactions with $\delta$-function interactions. The result is

$$R_V \sim \frac{2\pi V}{1+\pi^2 V^2},$$

$$R_J \sim 1 - \frac{1-\pi^2 V^2}{1+\pi^2 V^2} \frac{\ln(T/T_K)}{[\ln^2(T/T_K)+4\pi^2 S(S+1)]^{1/2}}.$$

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# Erratum

**Evaluation of the Partition Functions for Some Two-Dimensional Ferroelectric Models,** M. L. Glasser, [Phys. Rev. **284**, 359 (1969)].

(1) On the right-hand side of Eq. (13), $\lambda$ should be replaced by $\lambda^{-1}$.
(2) In Eq. (14) the argument of the logarithm should be

$$\Gamma(\alpha\beta/2\pi+\tfrac{3}{4})\Gamma(\alpha\gamma/2\pi+\tfrac{1}{4})/[\Gamma(\alpha\gamma/2\pi+\tfrac{3}{4})\Gamma(\alpha\beta/2\pi+\tfrac{1}{4})].$$

(3) To Eq. (23) add $= \ln|(2\mu/\pi)\cot(\pi^2/2\mu)\csc\mu|$.
(4) The right-hand side of Eq. (24) should read

$$(\tfrac{1}{8}\mu)I(\pi/2\mu, 3\mu, 5\mu).$$

(5) Equation (27) should read

$$z(0) = \ln|(2\mu/\pi)\cot\mu\cot(\pi^2/2\mu)|.$$

I wish to thank Dr. D. B. Abraham for pointing out the above simple form for Eq. (27).